

# Mechanical Energy = KE + APE

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①

Energy can be a useful way to understand fluid systems because all terms are expressed in the same units = Joules and Watts.

• Kinetic energy  $\equiv$  work done to accelerate a mass  $m$  to speed  $|\underline{u}|$ .

• Rate of doing work = force  $\times$  velocity (\*)

• In fluid systems we use  $\frac{\text{mass}}{\text{unit volume}}$   $\times$  acceleration =  $\frac{\text{force}}{\text{unit volume}}$

$$\text{or } \rho \frac{D\underline{u}}{Dt} = \underline{\text{Force}} / \text{vol.}$$

So to evaluate (\*) following a fluid parcel as it accelerates from 0 to speed  $|\underline{u}|$ :

$$\int_{t_0}^{t_1} \frac{\underline{\text{force}}}{\text{vol.}} \cdot \underline{u} dt = \rho \int_{t_0}^{t_1} \frac{D\underline{u}}{Dt} \cdot \underline{u} dt = \rho \int_{t_0}^{t_1} \frac{D}{Dt} \left( \frac{1}{2} \underline{u}^2 \right) dt$$

$$= \frac{1}{2} \rho \underline{u}^2 \quad \text{assuming } \underline{u} = 0 \text{ at } t = t_0$$

this is  $\frac{\text{KE}}{\text{unit vol.}}$   $\approx$  "KE<sub>v</sub>", units  $\frac{\text{J}}{\text{m}^3} = \frac{\text{kg m}^2}{\text{s}^2 \text{ m}^3} = \frac{\text{kg}}{\text{m s}^2}$

and often we casually (imprecisely) say  $\text{KE} = \frac{1}{2} \underline{u}^2 = \frac{m^2}{\text{s}^2}$

Note: our spatial integral for steady flow in the Bernoulli derivation gives the same result:

Notation  $\underline{x}^L$  = Lagrangian position of a fluid parcel

so each  $\underline{x}^L$  corresponds to a specific time

$\Rightarrow$  e.g. at  $t = t_0$  particle is at  $\underline{x} = \underline{x}_0^L$

$$\rho \int_{t_0}^{t_1} \frac{D\underline{u}}{Dt} \cdot \underline{u} dt = \rho \int_{t_0}^{t_1} \frac{D\underline{u}}{Dt} \cdot \frac{d\underline{x}^L}{dt} dt = \rho \int_{\underline{x}_0^L}^{\underline{x}_1^L} \frac{D\underline{u}}{Dt} \cdot d\underline{x}^L$$

$$\begin{matrix} \downarrow & & \downarrow \\ = \rho \frac{1}{2} \underline{u}^2 & \longleftarrow \text{same} \longrightarrow & = \frac{1}{2} \rho \underline{u}^2 \quad (\text{if } \frac{\partial}{\partial t} = 0) \end{matrix}$$

KE Derivation

Bernoulli derivation

Back to Energy:

For SW flow  $KE_v = \frac{1}{2} \rho (u^2 + v^2)$  and

we commonly use  $\frac{KE}{\text{unit horizontal area}} = "KE_A" = \int_{-h}^{\eta} KE_v dz$

$\Rightarrow KE_A = \frac{1}{2} \rho h (u^2 + v^2) \text{ for SW flow} \quad \text{Joules/m}^2$

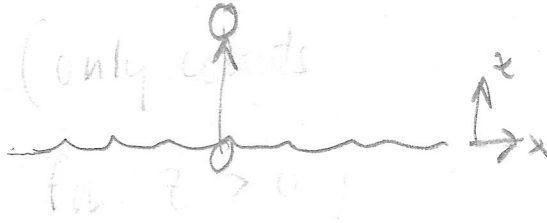
• Available Potential Energy  $\equiv$  Work done against gravity (buoyancy) to move a mass from its rest state to its current location.

Note: be sure you can define what the APE is available to do.

To move the free surface away from its rest state at  $z=0$ , for a single water particle

$$APE_v = \int_0^z \rho g dz = \rho g z$$

(only counts for  $z > 0$ )



only counts for above  $z=0$ . For downward displacements

To make the  $APE_A$  of a disturbed free surface we integrate over all such particles.

$$APE_A = \int_0^{\eta} \rho g z dz = \boxed{\frac{1}{2} \rho g \eta^2 = APE_A}$$

For downward displacement we are pushing air particles down through water - work done is still positive.

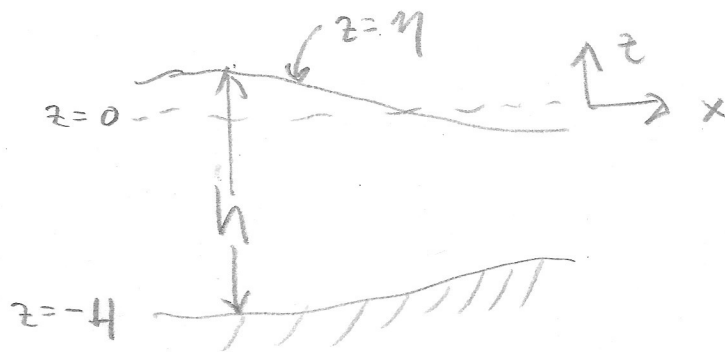
Forming governing equations for  $KE_A \circ APE_A$  (4)

- unidirectional, frictional, SW flow

- don't linearize advection terms (lead-in to hydraulics)  
 (lead-in to hydraulics)

$$\rho h u \cdot \left[ \boxed{x \text{ mom}} \quad u_t + u u_x + g \eta_x + \frac{C_d |u| u}{h} = 0 \right] \text{ (i)}$$

$$\rho g \eta \cdot \left[ \boxed{\text{mass}} \quad \eta_t + (h u)_x = 0 \right] \text{ (ii) note } \eta \eta_t = \left( \frac{1}{2} \eta^2 \right)_t \text{ eq.}$$



from (i):  $\rho h \left( \frac{1}{2} u^2 \right)_t + \rho h u \left( \frac{1}{2} u^2 \right)_x + \rho g h u \eta_x + \rho C_d |u|^3$

or:

$$\left( \rho h \frac{1}{2} u^2 \right)_t + \left( \rho h u \frac{1}{2} u^2 \right)_x - \left( \frac{1}{2} \rho u^2 \right) \left[ h_t + (h u)_x \right] \xrightarrow{0} \text{ by } \boxed{\text{mass}} \text{ note } h_t = \eta_t$$

$$+ (g \rho h u) \eta_x + \rho C_d |u|^3 = 0$$

from (ii)

$$\left( \frac{1}{2} \rho g \eta^2 \right)_t + \rho g \eta (h u)_x = 0$$

these combine to form "pressure work"

Adding the results from (i) + (ii)  
 we find the full SW energy equation

$$\underbrace{\left( \frac{1}{2} \rho h u^2 + \frac{1}{2} \rho g \eta^2 \right)_t}_{\text{Rate of change of KE}_A + \text{APE}_A} + \underbrace{\left[ u \left( \frac{1}{2} \rho h u^2 \right) + \rho g \eta h u \right]_x}_{\substack{\text{(Divergence of)} \\ \text{Advection of KE}_A}} = \underbrace{-\rho C_d |u|^3}_{\substack{\text{(Divergence of)} \\ \text{Pressure work} \\ \star}}$$

Energy loss to bottom friction (neg. definite)

★ Pressure work =  $(\rho g \eta) \times (hu)$  = pressure anomaly × transport  
 - is how energy is transmitted over long distances by waves

- Except for very fast flow, Pressure Work is much larger than KE Advection:

Scaling  $\frac{\text{KE Adv.}}{\text{Press. Work}} = \frac{\frac{1}{2} u^3 H}{\rho g H u} \sim \frac{u^2}{g \left( C \frac{a}{H} \right)} \sim \frac{u}{C}$

This is called a "Froude Number" ~ u from linear SW waves  
 (velocity / wave speed)

For our linear SW waves;

(6)

$$Fr = \text{Froude \#} \sim \frac{u}{c} \sim \frac{a}{H} \ll 1$$

$\Rightarrow$  KE advection unimportant

Also:

- $c_g = c_p = \sqrt{gH}$  : non-dispersive

- Scales of  $KE_A$  &  $APE_A$  are equal ("equipartition")

(★) Prove this, using  $u \sim c \frac{a}{H}$

- Can estimate the net energy flux through a channel of width  $B$  as

$$\text{Flux} \sim B \rho g H \langle \eta u \rangle \quad \langle \rangle = \text{tidal average}$$

(★) • What are the units of the flux?

- What if  $\eta$  &  $u$  are out of phase?

- What does this tell you about net energy dissipation (say channel leads to a closed bay)?